Bayesian Combined Analysis of JET LIDAR, Edge LIDAR and Interferometry Diagnostics
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1. INTRODUCTION
The accurate determination of electron density $n_e$ and temperature $T_e$ is important for many areas of Tokamak research. At JET, several different systems make measurements which depend on these but are based on very different physics, have different spatial resolutions and observe different regions of the plasma. Bayesian analysis and forward modelling provide a conceptually simple method to obtain and represent the information that can be inferred from the combination of an arbitrary number of observations and prior assumptions about all the physical quantities on which they depend. This paper gives results from the Bayesian Inversion for a typical JET plasma using models developed for the Interferometry [1], core LIDAR [2] and edge LIDAR [3] diagnostics.

2. MODELS
For each diagnostic, a ‘forward model’ is created which can predict the distribution of measurements which might be observed given a predefined exact physical state. The interferometry is modelled simply by integrating $n_e$ along its lines of sight but the LIDAR model is considerably more complex, depending on many extra calibration/auxiliary parameters. Figure 1 shows the lines of sights of the three diagnostics and an overview and sample output of the LIDAR model.

In the standard analyses, the auxiliary parameters must be fixed to calibration measurements or determined by statistical cross calibration with other diagnostics. While these can be used as the prior distribution for the parameters here, there is sufficient information from the combined systems to use weak priors in many cases. In particular, because the spectrometer sensitivities of the edge LIDAR system are difficult to determine (and to demonstrate the applicability of the technique to cases where some calibration parameters are completely unknown), these parameters are assigned effectively no prior information - a uniform distribution. The timing parameters which effect inferred positions, are also assigned uniform priors.

For the plasma, $n_e$ and $T_e$ are assumed constant on any poloidal flux surface so are modelled as 1D functions of poloidal flux $\psi_N$ (normalised to 0 at the magnetic axis and 1 at the last closed flux surface), which is taken from the JET routine equilibrium code [4]. A linear interpolation of 40 ‘knots’ at fixed positions is used, with more in the H-Mode pedestal region (see figure 2a/b). The knot magnitudes are the parameters for which the posterior is obtained.

3. RESULTS - A TYPICAL POSTERIOR
The posterior distribution gives the probability for any combination of the plasma and calibration parameters and so describes all of the information that can be known about these parameters, given all the prior assumptions and all the observations. It includes all uncertainties, systematic and random, from all modelled sources. (For previous examples of the procedure, see e.g. [5, 6]). To examine this high-dimensional distribution, a series of representative samples of are drawn. Each sample is a complete description of a possible state of the entire system, that is consistent with all the diagnostics. They can be displayed separately or used to generate histograms for each parameter, giving the marginal distribution which expresses what can be inferred about that parameter, independent of all others. Figure 2 shows several samples of the $n_e$ and $T_e$ profiles for a time point in a typical H-
Mode JET plasma, as well as the marginal distributions for regular points along the \( N_e \) and \( T_e \) profiles from the same posterior. The time point used lies between ELMs, where the profiles are unlikely to have evolved over the few milliseconds between the capture/integration times of different diagnostics.

Despite the complete freedom (weak prior distributions) given to many of the calibration parameters which must be fixed in the standard approach, the results are good and many benefits of the integrated approach can be seen. The overall magnitude of the density differs to the standard analysis of core LIDAR due to the inclusion of the Interferometry diagnostic’s very accurate integral information. The profiles also show much more is inferred in the pedestal region than the core LIDAR standard analysis shows. While this appears entirely due to the edge LIDAR data, core LIDAR provides much of the information. Because the priors given for many of the edge LIDAR calibration parameters were very weak in this case, edge LIDAR alone is not sufficient to determine the pedestal profiles. For instance, while the edge LIDAR data can be used to find the ne pedestal shape, it cannot give its magnitude as the overall sensitivity of the optics is not known and so the prior on that parameter was uniform. The same is true for core LIDAR, meaning it provides the edge pedestal density only relative to the core. The interferometry absolute density information completes the picture and so, without stronger prior information, only together can the three systems provide the profile.

The \( T_e \) pedestal is similar because the priors for the edge LIDAR spectrometer sensitivities were uniform. The edge LIDAR data alone does heavily constrain the possible sensitivity combinations because the same set must be consistent for the entire profile but there remains a degree of freedom which is always enough to allow any pedestal \( T_e \). The core LIDAR system, with well determined sensitivities, can in principal provide the edge temperature but lacks the resolution to do this by itself. Consideration of the instrument convolution (which is included in the model but not in the standard analysis) increases the effective resolution but shows that practically, it is a ~ 12cm spatial integral of \( T_e \) that is really known. Combined with shape information from the edge LIDAR data, it is just sufficient to reconstruct the profile. Figure 3 shows the inferred \( T_e \) profile for just the edge region plotted against major radius. Also shown is the result from the standard analysis of the edge LIDAR data, using values for the spectrometer sensitivities which have not been validated. The apparent discrepancy is simply due to the uncertainty in these parameters which, while difficult to identify and treat when interpreting the results, is rigorously and inherently handled with the Bayesin approach.

The posterior distribution also describes what can be inferred about the calibration parameters themselves. For a selection of times in 2 pulses, samples of the posterior obtained and the ratios of the spectrometer sensitivities taken from each (shown in Figure 4). Each one represents what can be known about that parameter from the model, prior and that data. Shown as bands is the combined distribution which could be used as a prior for future inversions.

4. PEDESTAL PARAMETERISATION

It is often useful to use a stronger plasma parameterisation in order to directly obtain the posterior over quantities of interest. For instance, the \( n_e \) pedestal is believed to be well approximated by a hyperbolic tangent function[9] of which the width is of particular interest as its scaling with parameters such as ion gyro radius have particular importance for ITER [10]. Figure 5 shows the posterior
distribution for the pedestal width \( w \), obtained using the parameterisation \( n_e = n_e^{\text{ped}} \left( 1 + \tanh(2(\psi_0 \psi_N) = w) \right) \) added to the knot parameterisation for the core. Samples of \( n_e(\psi_N) \) are also shown. The strong assumption of shape, when valid, is particularly useful if the LIDAR signals are weak as in this case where the data has a noise level of \( \sim 15\% \), which can be seen to strongly affect the standard analysis.

In future work, the posterior will be found for a series of pulses where the quantities of interest (e.g. \( \rho \)) were varied. From these the information that can be inferred from the LIDAR systems about any relationship of the pedestal quantities on those parameters will be found, supplementing the analysis using the High Resolution Thomson Scattering[10].

CONCLUSIONS
The development of a model of the LIDAR Thomson Scattering diagnostic had been outlined and its use to infer electron density and temperature profiles consistent with three systems at JET has been demonstrated. The detail of the model allows the uncertainties and complex relationships of various calibration parameters to be handled easily and rigorously. It is now necessary to include a model for the plasma current and magnetic diagnostics (from [11]) in order to include in the uncertainties, those from \( \psi_N \), to which the combination of the core and edge LIDAR systems is very sensitive, due to their distant lines of sight. The authors would like to thank A. Meakins for work on the Genetic Algorithm used.

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REFERENCES
Figure 1: a) Poloidal cross-section of JET and lines of sight of interferometry (dotted), core and edge LIDAR (dashed) and typical flux surface geometry (gray). b) LIDAR Forward model outline using physics parameters (yellow) and calibration/other parameters (light yellow). c) Core LIDAR digitiser trace for spectrometer channel 2 showing observed data (blue) and likelihood distribution mean (red) and ±2σ (light red) from the model.

Figure 2: Profiles of electron density (a) and temperature (b) vs normalised poloidal flux of 8 samples from the posterior distribution inferred from core LIDAR, edge LIDAR and interferometry with weak calibration parameter priors. c/d) Marginalised posterior distribution for points along normalised poloidal flux. Also shown are the parametrisation knot positions (vertical lines on a/b) and profiles from routine analysis of the High-Resolution Thomson Scattering[7] (orange ∆), Heterodyne Radiometer Electron Cyclotron Emission[8] (green ◊) diagnostics and the standard analysis of core LIDAR (magenta * and line).
Figure 3: $T_e$ profile samples (left) and marginal distributions (right) on the magnetic axis plane ($Z = Z_{mag}$) versus major radius $R_{mag}$ for the outboard plasma edge. Other diagnostics as in figure 2 with edge LIDAR standard analysis (black/white +) based on unvalidated spectrometer sensitivity values.

Figure 4: Posterior distributions for sensitivity ratio of two edge LIDAR spectrometers for several time points (2 pulses). Bars show mean and $1\sigma$ (black), $2\sigma$ (dark gray) and $3\sigma$ (light gray). Bands show combined result.
Figure 5: a) Posterior sample profiles (blue lines) determined from tanh parameterisation and others as in figure 3, b) Marginal posterior PDF for pedestal width.